



# Renormalization and UV behaviour of 5D gauge theories

Anca Preda

in collaboration with G. Cacciapaglia, W. Isnard, R. Pasechnik



**LUND**  
UNIVERSITY

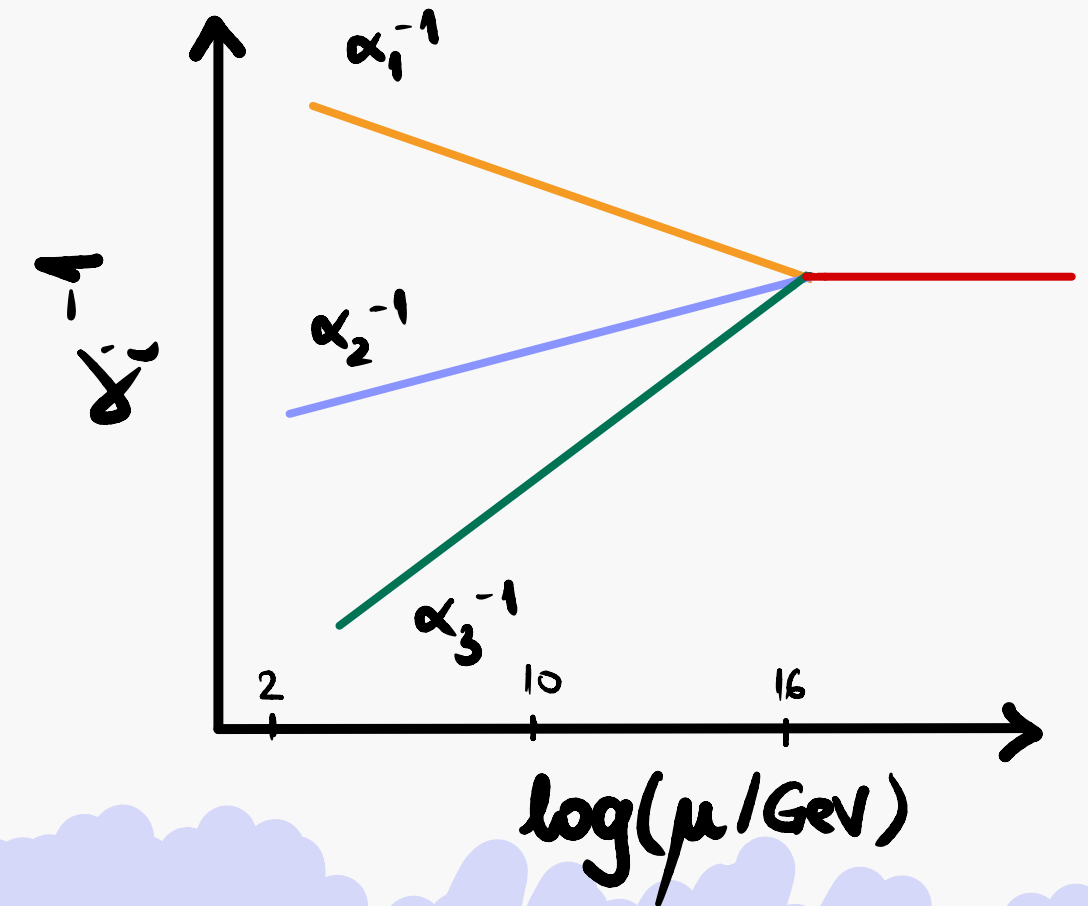
# Outline

- Motivation
- Running of couplings (RGEs): from 4D to 5D
- Renormalization of 5D gauge theories
- Conclusions



# Asymptotic Grand Unification Theories (GUTs)

# Asymptotic Grand Unification Theories (GUTs)



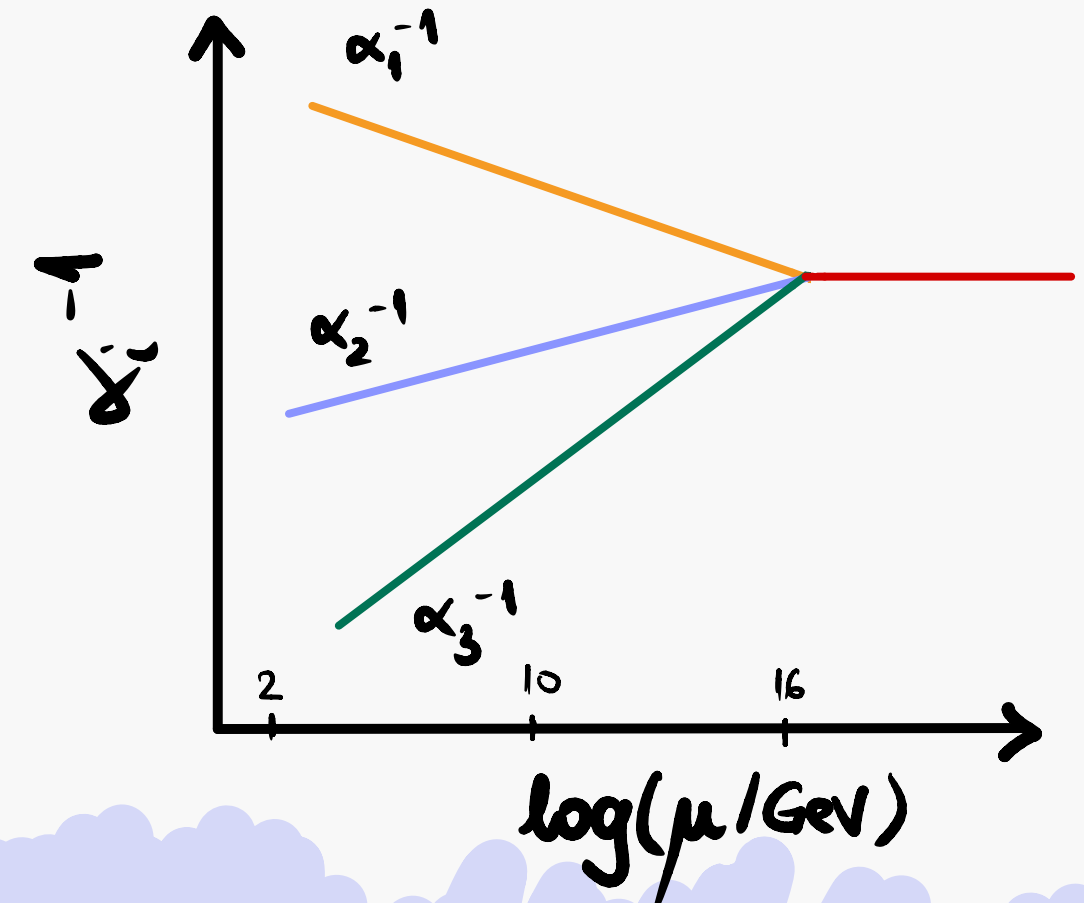
4D theories

standard idea of gauge coupling unification

SU(5), SO(10) ...

# Asymptotic Grand Unification Theories (GUTs)

extra dimensional theories (5D)  
same models SU(5), SO(10) ...  
different construction



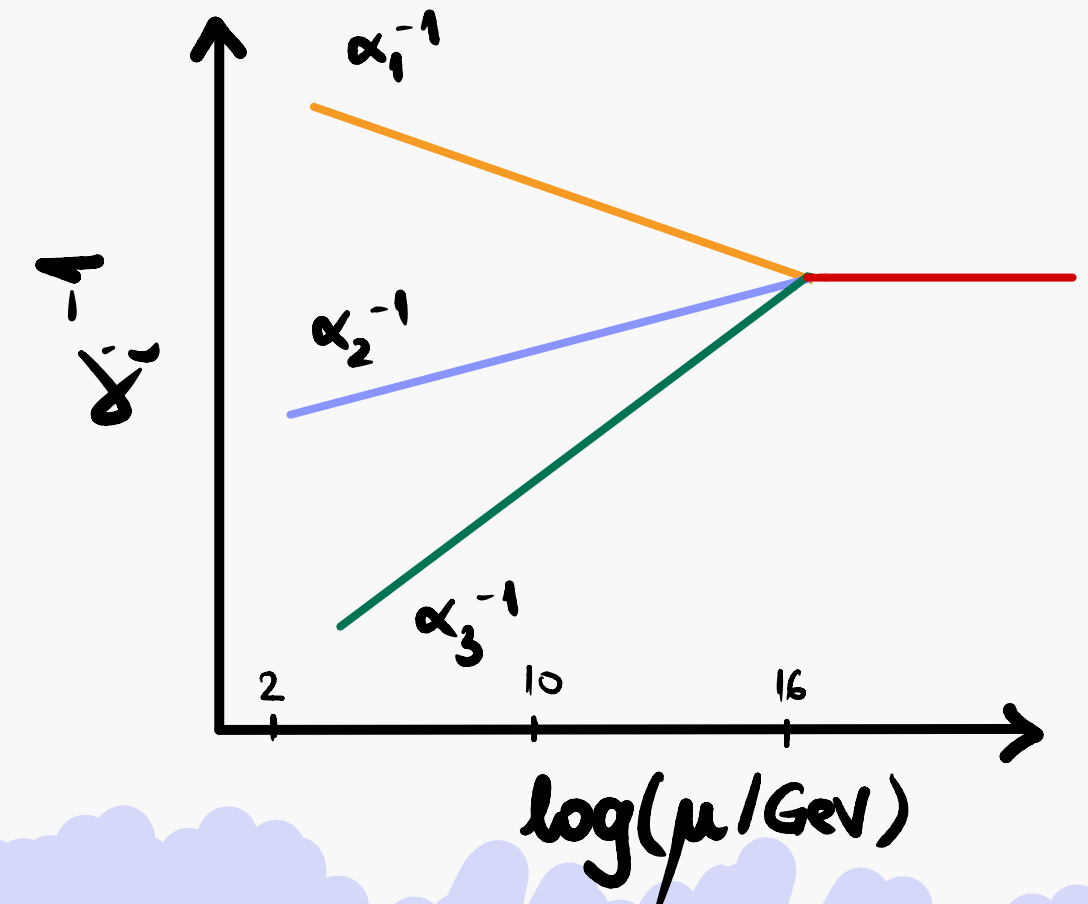
4D theories  
standard idea of gauge coupling unification  
SU(5), SO(10) ...

# Asymptotic Grand Unification Theories (GUTs)

extra dimensional theories (5D)  
same models SU(5), SO(10) ...  
different construction

## Why?

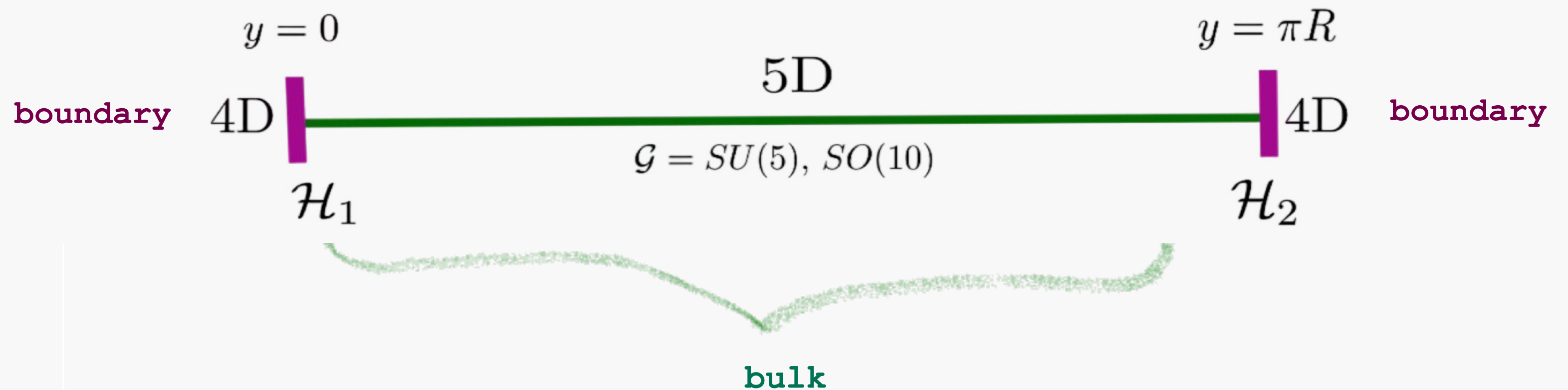
- lower unification scale
- less parameters
- hierarchy problem



4D theories  
standard idea of gauge coupling unification  
SU(5), SO(10) ...

# Asymptotic GUTs

- **5D models:** one extra dimension<sup>1, 2</sup> compactified on  $K = S^1/\mathbb{Z}_2 \times \mathbb{Z}'_2$



<sup>1</sup> A. Hebecker, J. March-Russell, Nuclear Phys. B 625 (2002)

<sup>2</sup> G. Cacciapaglia, arXiv:2309.10098 (2023)

Any 5D field

$$\Phi_{5D}$$

(scalar, fermion, gauge boson)

Infinite tower of 4D fields

⋮

n=4 —————

n=3 —————

n=2 —————

n=1 —————

n=0 —————



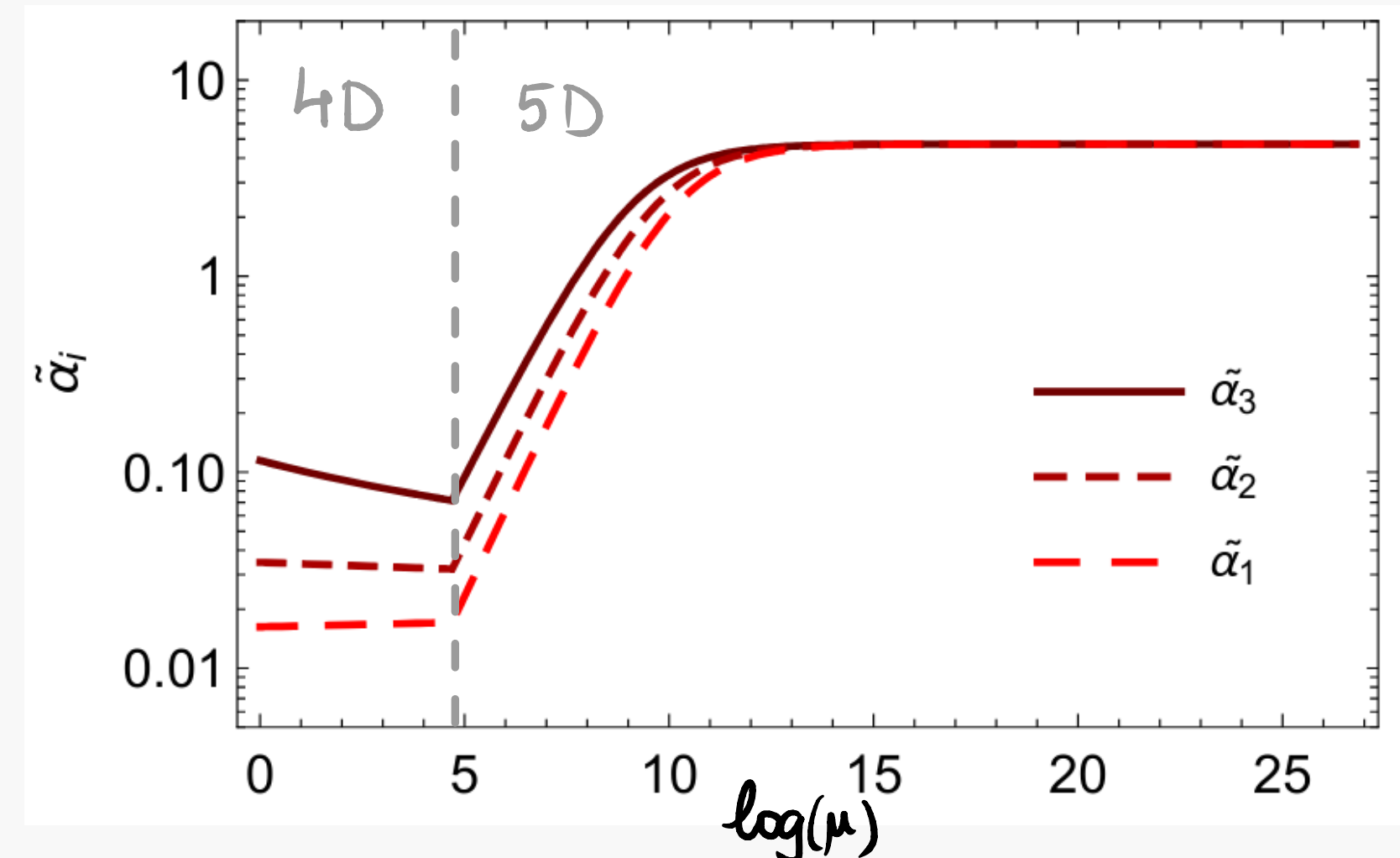
# RGEs: from 4D to 5D

- 4D: logarithmic running  $16\pi^2 \frac{dg_i}{dt} = b_{SM}^i g_i^3$

- 5D: power-law running <sup>2</sup>

$$16\pi^2 \frac{dg_i}{dt} = b_{SM}^i g_i^3 + (S(t) - 1)b_5^i g_i^3$$

$$\text{where } S(t) = \begin{cases} \mu R = M_Z R e^t, & \text{for } \mu \geq 1/R \\ 1, & \text{otherwise} \end{cases}$$



G. Cacciapaglia, et.al., 104, PRD 2021

<sup>2</sup> G. Cacciapaglia, arXiv:2309.10098 (2023)

# RGEs: from 4D to 5D

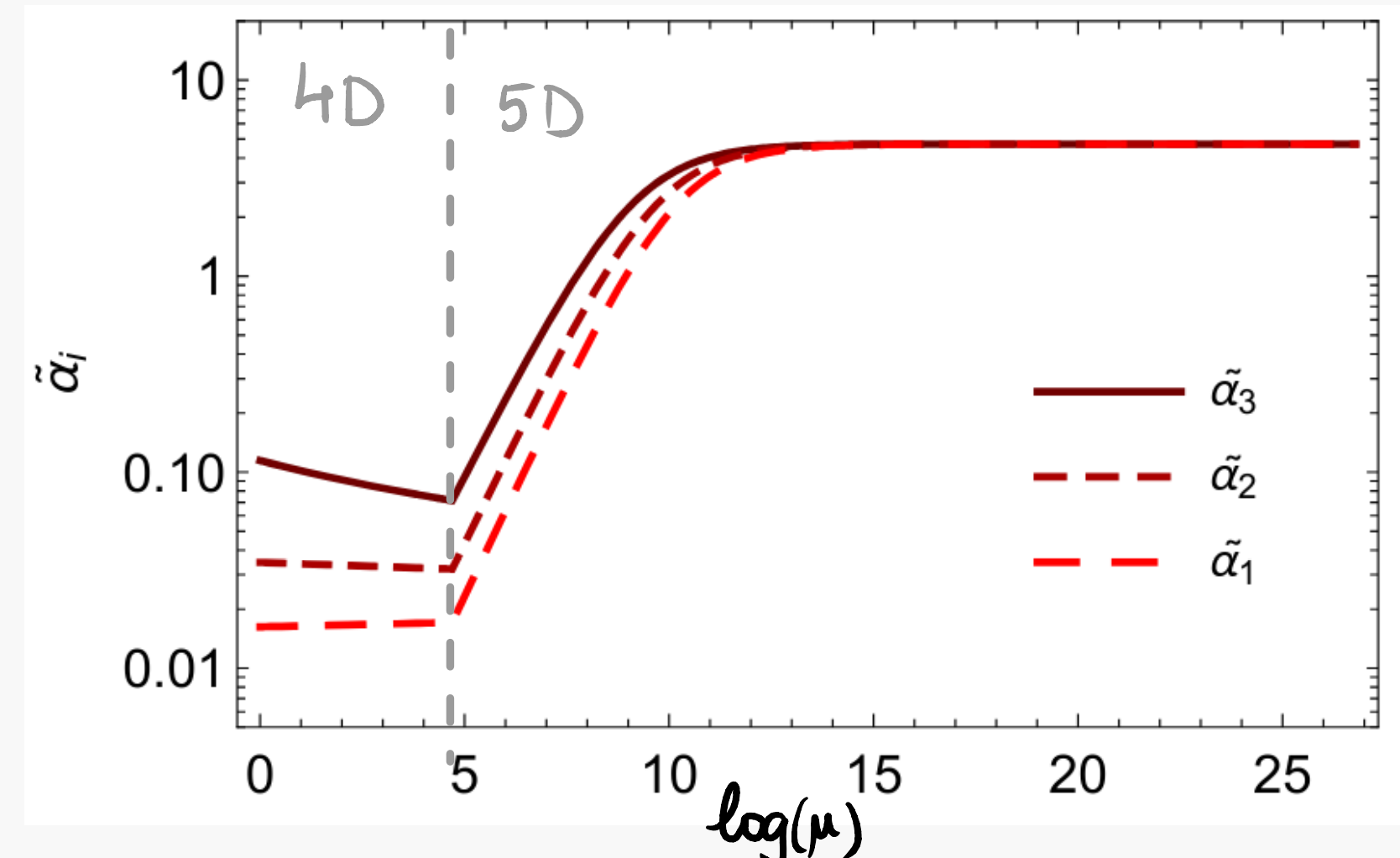
- 5D: couplings flow asymptotically towards a UV fixed point

It's existence  $\longleftrightarrow$  good behavior in the UV

- gauge couplings

$$16\pi^2 \frac{dg_i}{dt} = b_{SM}^i g_i^3 + (S(t) - 1)b_5^i g_i^3$$

Fixed point exists for  $b_5^i \leq 0$



G. Cacciapaglia, et.al., 104, PRD 2021

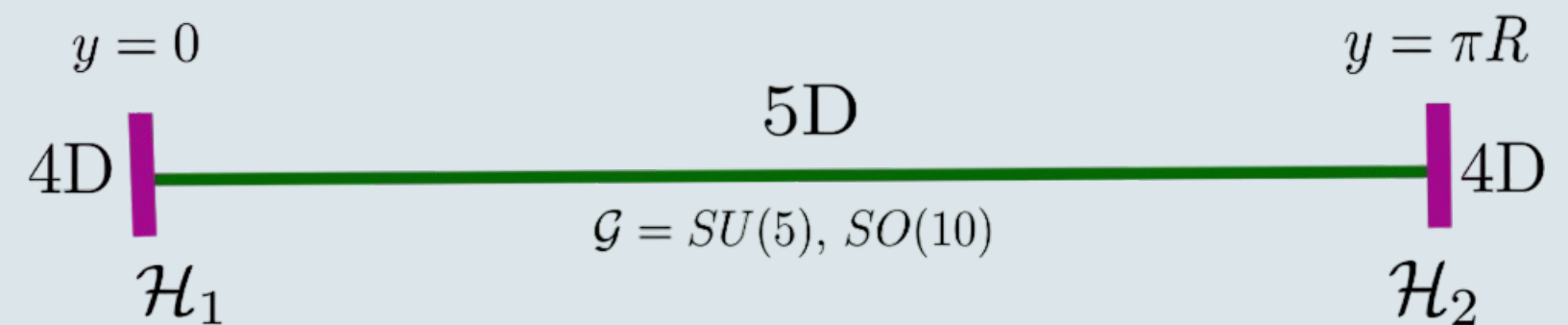
---

Assume a **theory**  
with couplings that  
flow asymptotically  
towards **fixed points**

+ the fixed points are perturbative

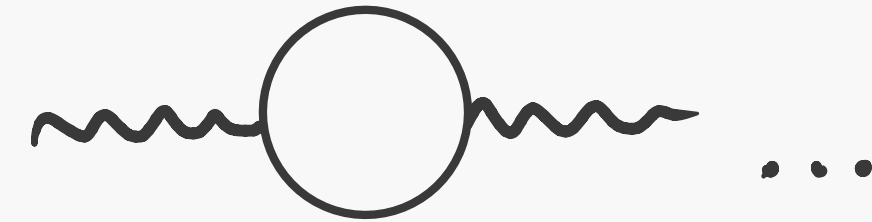
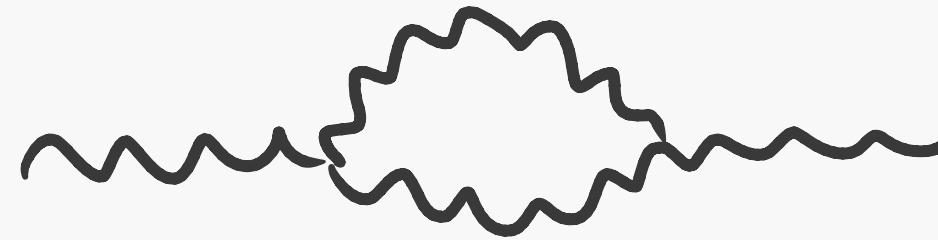


Can we **renormalize**  
the theory in the bulk  
and on the  
boundary?



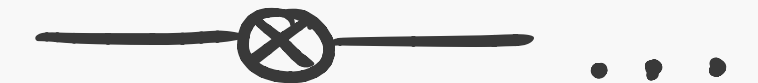
# Renormalization

Loop diagrams may be divergent  $(+\infty)$



Introduce counterterms to cancel divergencies

$(-\infty)$



If the number of counterterms is finite: theory is **renormalizable**



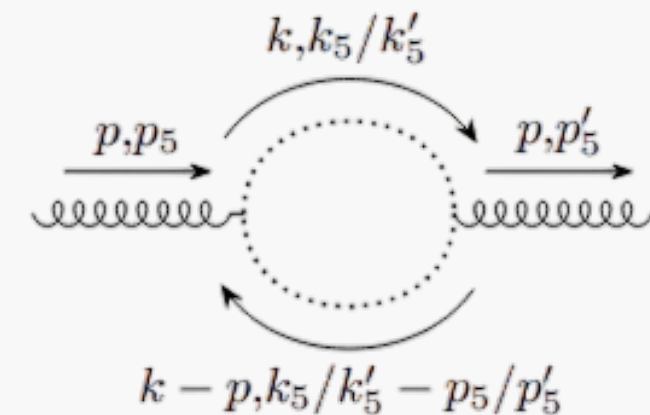
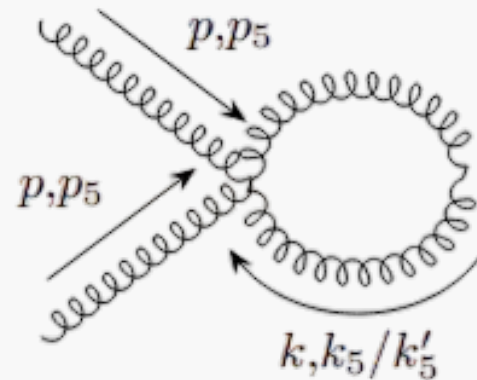
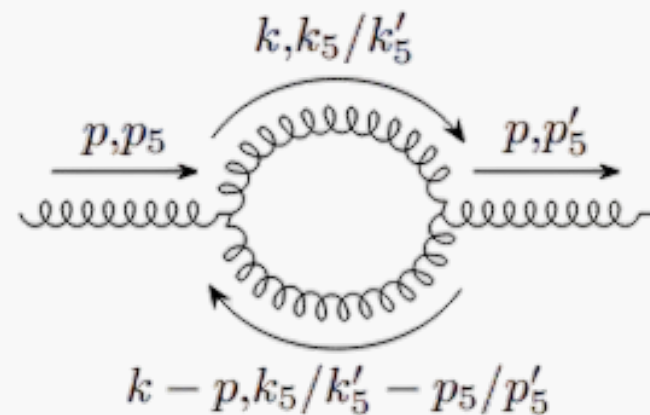
$(\text{finite})$

# Renormalization of 5D gauge theories

Common lore: **5D** gauge theories are **non-renormalizable** <sup>3</sup>

... but under certain conditions a fixed point exists  $\rightarrow$  theories valid up to arbitrarily high energy scales <sup>4</sup>

Pure Yang Mills theory  $\mathcal{L}_{\text{YM}} = -\frac{1}{4}F^{MN}F_{MN}$



<sup>3</sup>H. Gies, Phys. Rev. D 68 (2003)

<sup>4</sup>T. Morris, JHEP 2005 (2005)

# Renormalization of 5D gauge theories

Pure Yang Mills theory  $\mathcal{L}_{\text{YM}} = -\frac{1}{4}F^{MN}F_{MN}$

## Bulk: 5D Yang Mills

- structure of divergencies same as in 4D
- scaling changes from 4D log to 5D linear
- **finite number of counterterms**

## Boundary: Yang Mills at the fixed points

- **finite number of counterterms**
- “magic gauge” ( $\xi = -3$ ): single counterterm  
$$-\frac{1}{4}K\delta(F_{\mu\nu}F^{\mu\nu})$$

**! Theory is renormalizable both in the bulk and on the boundary (at one loop)**

( finite number of operators to cancel divergencies )

# Renormalization of 5D gauge theories

## Adding scalars and fermions

- realistic models: go beyond pure Yang Mills
- **Yukawa interactions**: finite on the boundary<sup>5</sup>
- fermion and scalar 2-pt functions introduce finite number of counterterms<sup>6</sup>
- **scalar quartic interactions** also renormalizable

<sup>5</sup>H. Georgi, et al, Phys. Let. B 506 (2001)

<sup>6</sup> HC Cheng, et al., Phys. Rev. D, 66 (2002)



---

## Summary and conclusions

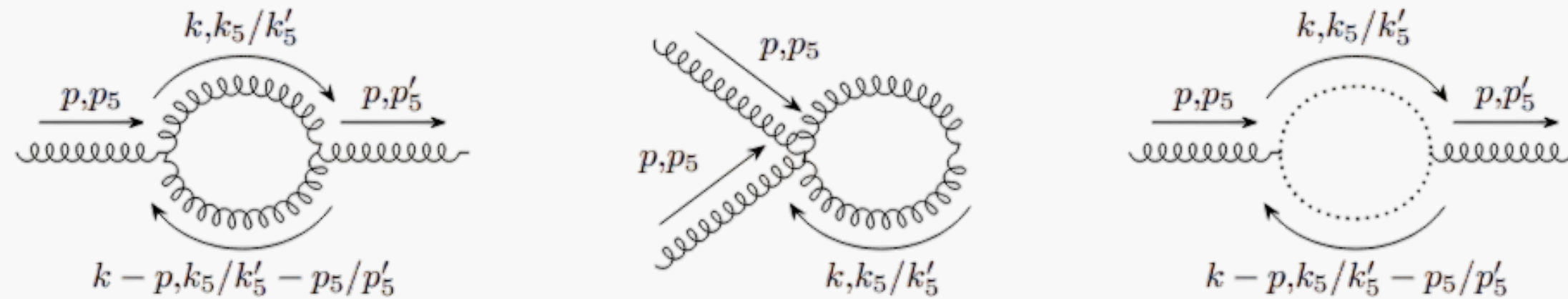
- RGEs of couplings change in **5D: power law running** (compared to 4D logarithmic)
- Consistency of models requires the existence of **UV fixed points**
- Fixed points  $\longleftrightarrow$  renormalizability
- Pure **Yang Mills renormalizable** at one loop (both bulk and boundary)
- Going beyond Yang Mills: also renormalizable



**Backup slides**

# Renormalization of 5D gauge theories

Pure Yang Mills theory  $\mathcal{L}_{\text{YM}} = -\frac{1}{4}F^{MN}F_{MN}$



$$i\Sigma = [p_5 - \text{conserving terms}] + [p_5 - \text{non-conserving terms}]$$

**Bulk:** structure of divergencies same as in 4D (at one loop)

- Scaling changes from logarithmic 4D ( $\sim \log \Lambda$ ) to 5D linear ( $\sim \Lambda$ )
- Same renormalization procedure as in 4D; **finite number of counterterms** are needed

# Renormalization of 5D gauge theories

Pure Yang Mills theory  $\mathcal{L}_{\text{YM}} = -\frac{1}{4}F^{MN}F_{MN}$

**Boundary:** renormalizability of a pure Yang Mills at the fixed points

→ less straightforward than bulk

Two-point function: <sup>3,6</sup>

$$i\Sigma = \frac{g^2}{64\pi^2} \frac{1}{\epsilon} \left[ (g_{\mu\nu} - p_\mu p_\nu) \left( \frac{11}{3} - (\xi - 1) \right) C(G) + g_{\mu\nu} \frac{p_5^2 + p_5'^2}{2} (4 + (\xi - 1)) C(G) \right]$$

To reconstruct the full counterterm on the boundary we need **higher vertex corrections** as well (3-pt and 4-pt functions)

(work in progress)

<sup>3</sup>G. Cacciapaglia, W. Isnard, R. Pasechnik, **AP**, (in preparation)

<sup>6</sup> HC Cheng, et al., Phys. Rev. D, 66 (2002)

# Renormalization of 5D gauge theories

Pure Yang Mills theory  $\mathcal{L}_{\text{YM}} = -\frac{1}{4}F^{MN}F_{MN}$

**Boundary:** renormalizability of a pure Yang Mills at the fixed points

- less straightforward than bulk
- divergencies can be absorbed by a **finite number of counterterms**
- “**magic gauge**” (  $\xi = -3$  ): coefficients of the 2-,3- and 4-pt functions are the same  
localized counterterm can be unified into a single term  $-\frac{1}{4}K\delta(F_{\mu\nu}F^{\mu\nu})$

**! Theory is renormalizable both in the bulk and on the boundary (at one loop)**

( finite number of operators to cancel divergencies )