The background of the slide is a vibrant cosmic scene. It features a large, bright yellow and orange galaxy core on the left, surrounded by blue and purple nebulae. The right side shows a dense field of stars and a bright, glowing region. The overall color palette is dominated by deep blues, purples, and oranges, creating a sense of depth and vastness in space.

Stochastic gravitational waves from supercooled cosmological phase transitions

Roman Pasechnik
Lund University

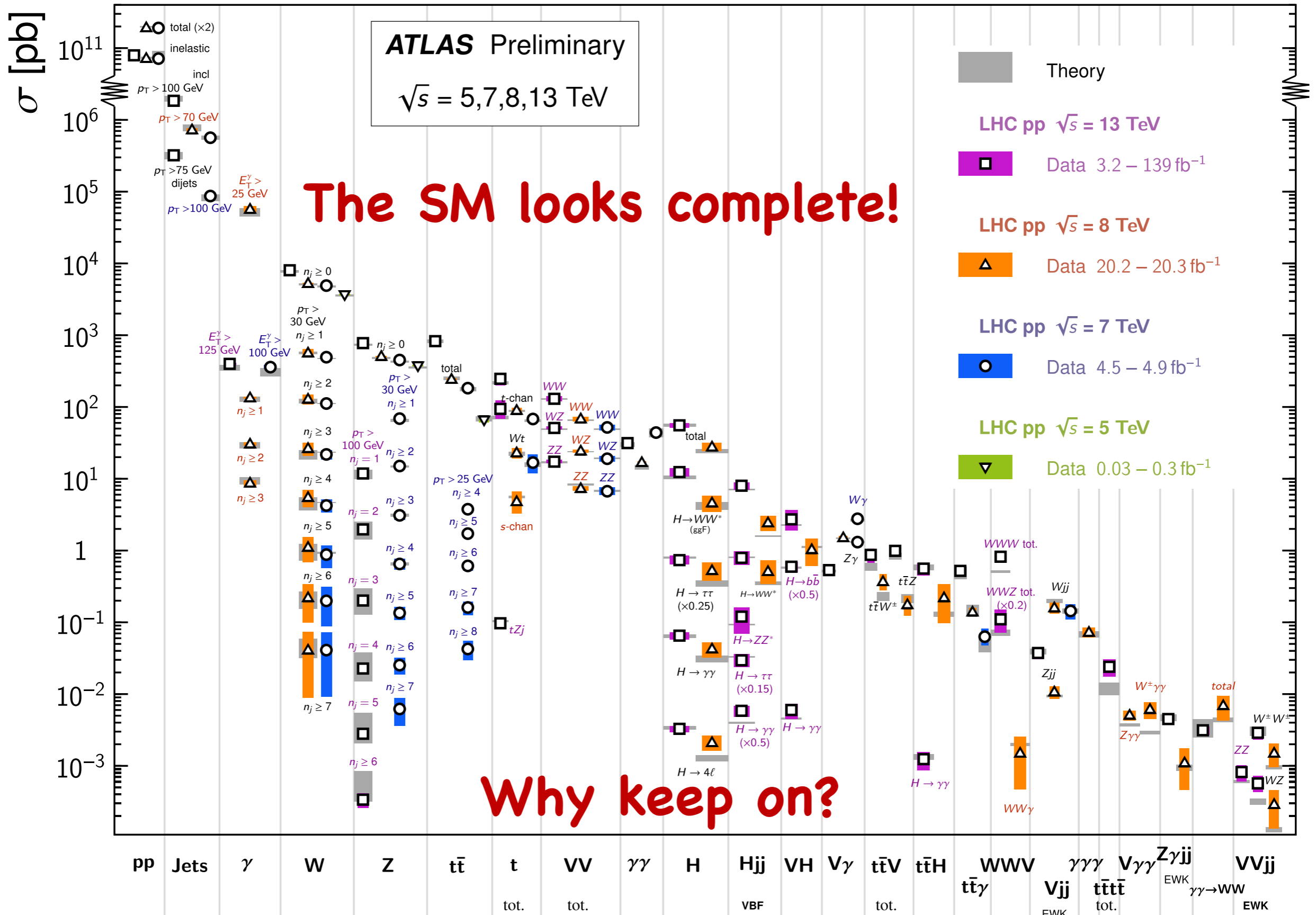
In collaboration with:
João Gonçalves, Antonio Morais, Danny Marfatia

2nd National SKA Science Day Sweden

Impressive performance of the Standard Model

Standard Model Production Cross Section Measurements

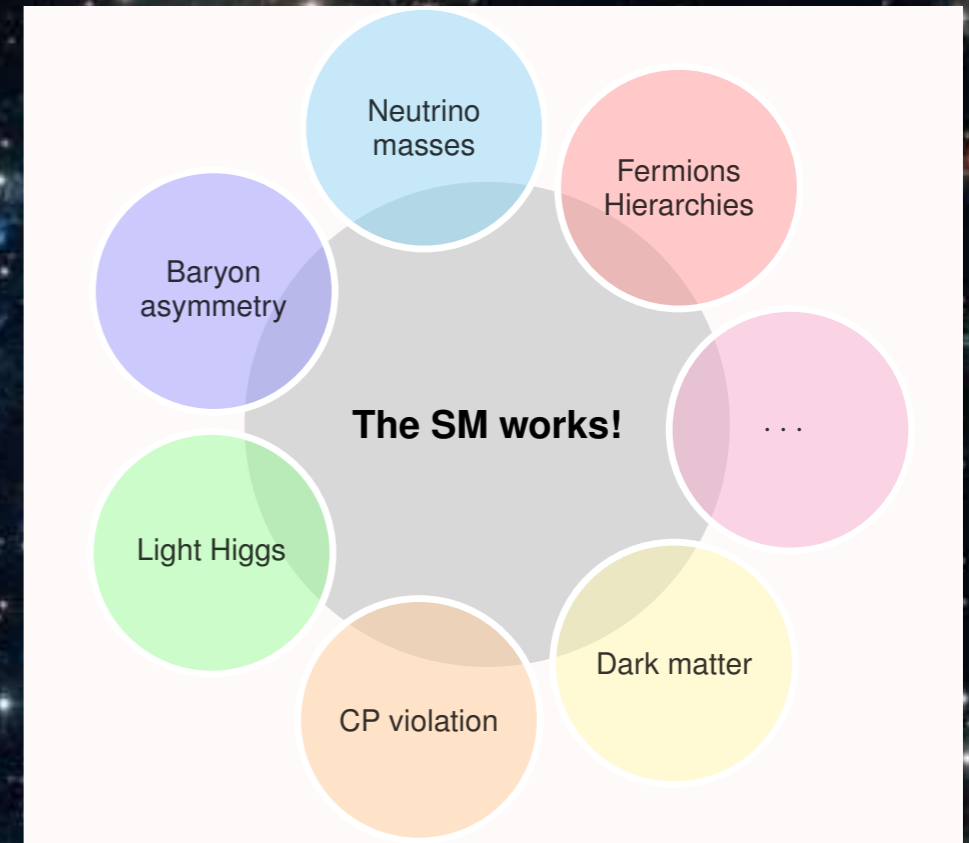
Status: February 2022



Open questions

Unexplained phenomena:

- Dark Matter
- Matter-antimatter asymmetry
- Dark Energy



Unsatisfactory structure of the SM:

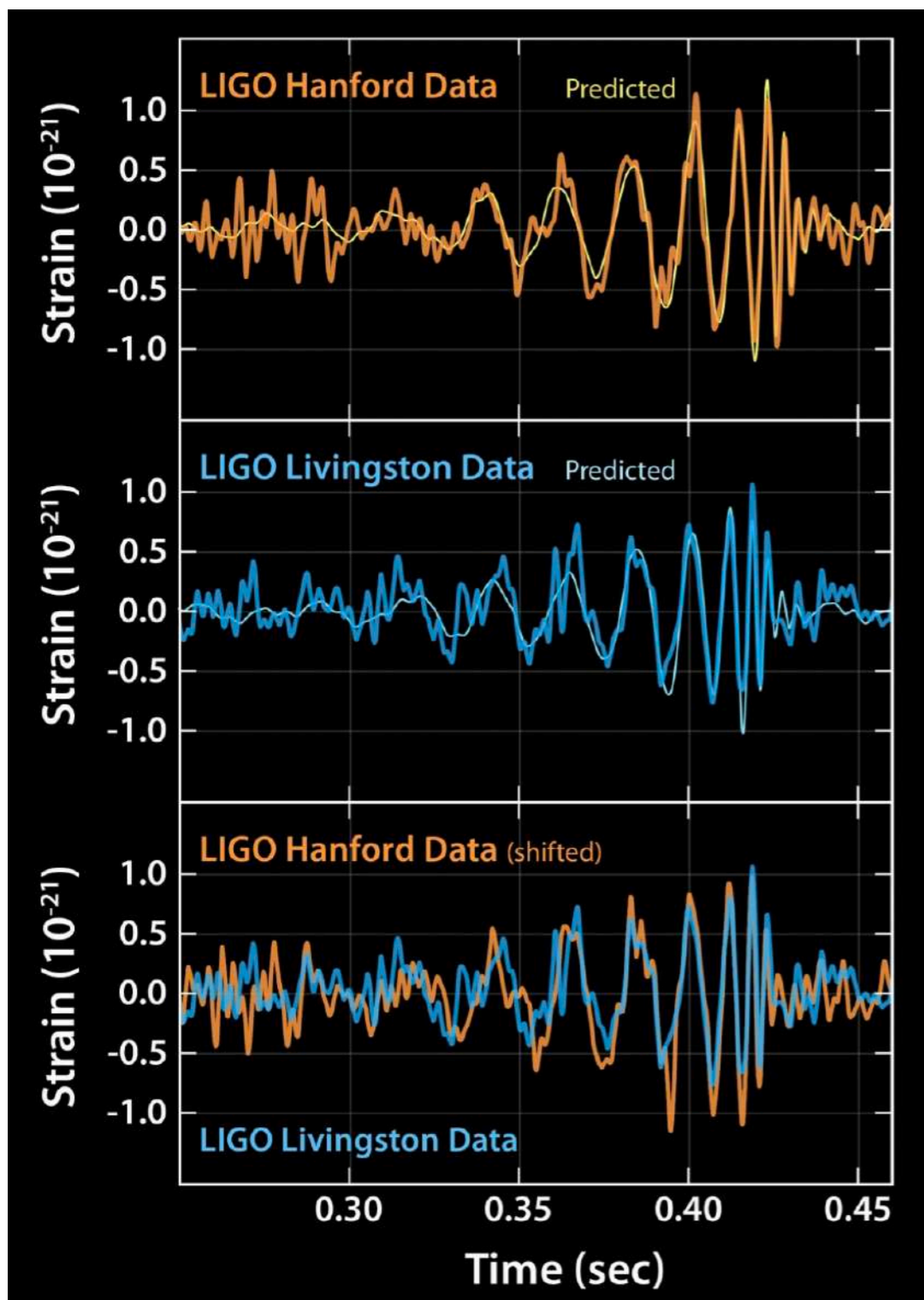
- Hierarchy problems (Higgs, flavour)
- Naturalness
- Quantum Gravity etc

What are the alternative ways to probe New Physics?

New Era of gravitational-wave astrophysics

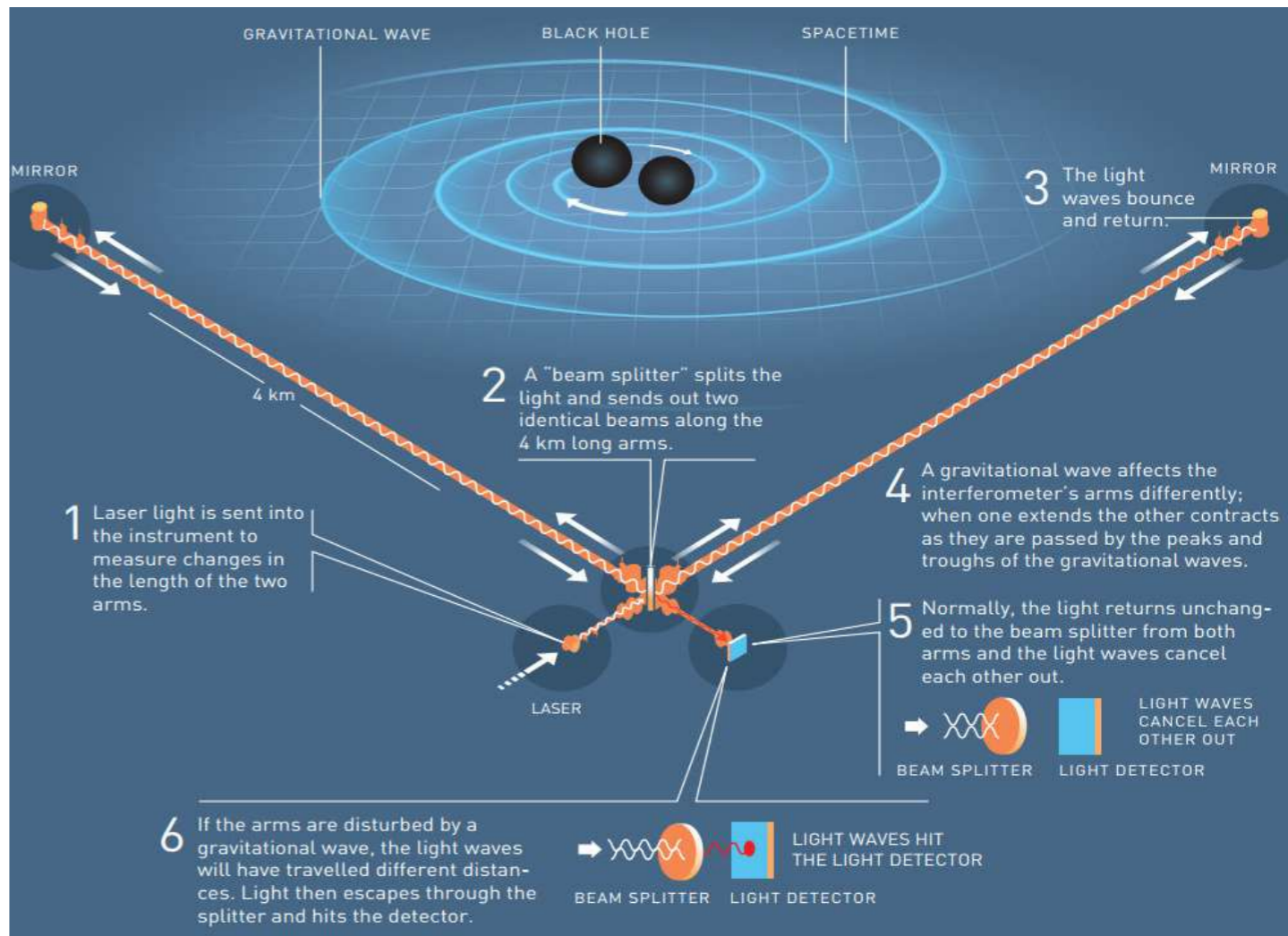
The first ever detection of gravitational waves

Event **GW150914**



Credit: LIGO

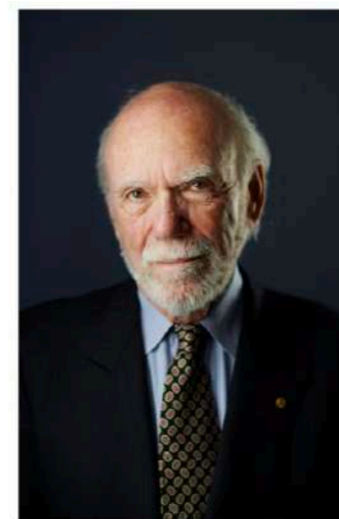
LIGO spectrometer



The Nobel Prize in Physics 2017



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Barry C. Barish



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Kip S. Thorne

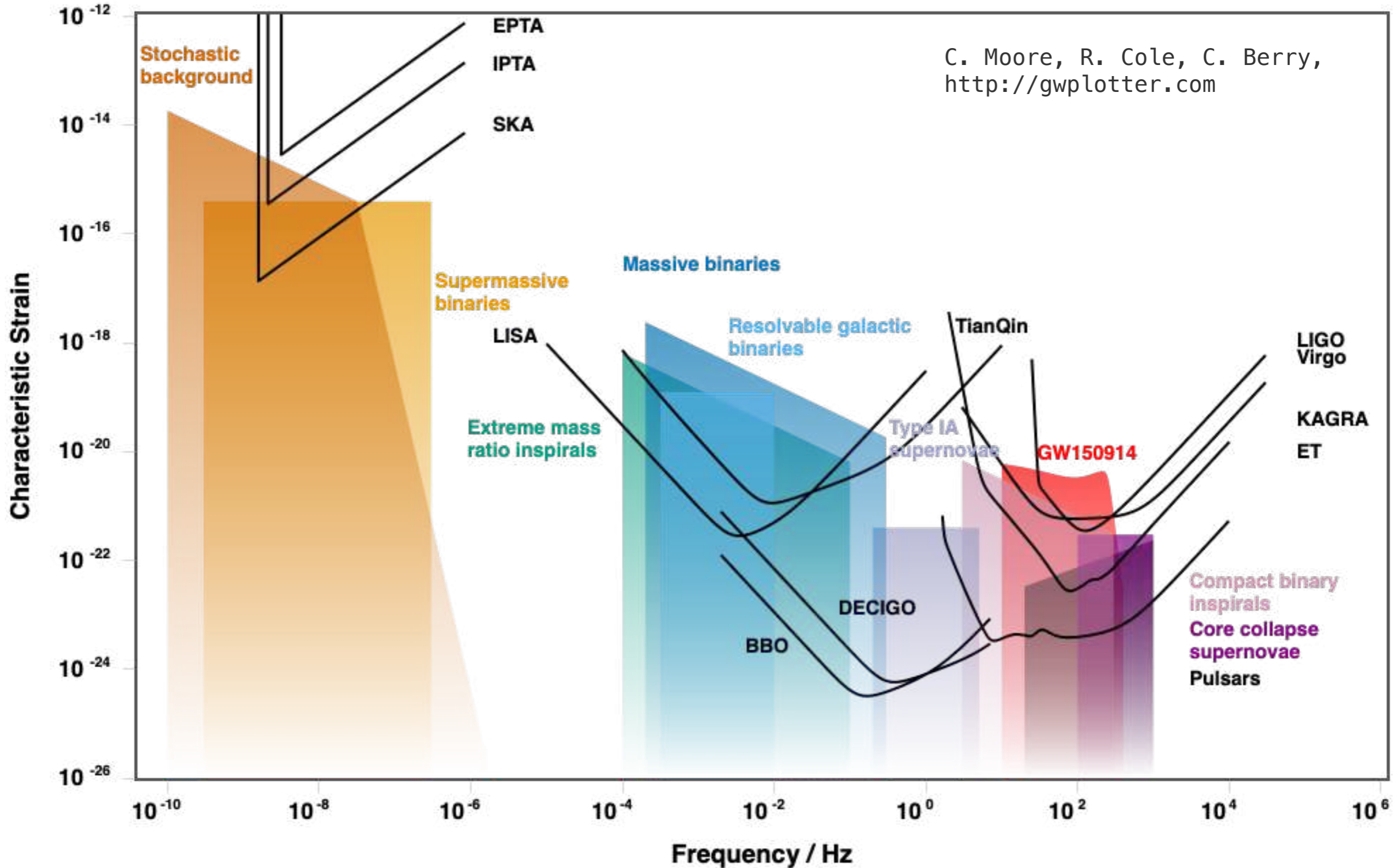
Why do Gravitational Waves matter?

Quick answer: GWs, in the form of a stochastic cosmological background, allow to probe physics not at the reach of collider experiments .

- Cosmological events

- (i) Inflation
- (ii) Cosmic strings
- (iii) **Strong cosmological phase transitions (PTs)** → by expanding vacuum bubbles of a broken phase in a universe filled with a symmetric phase

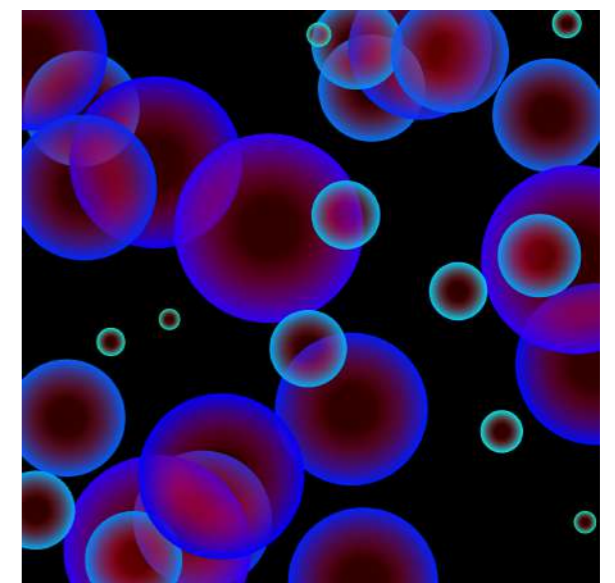
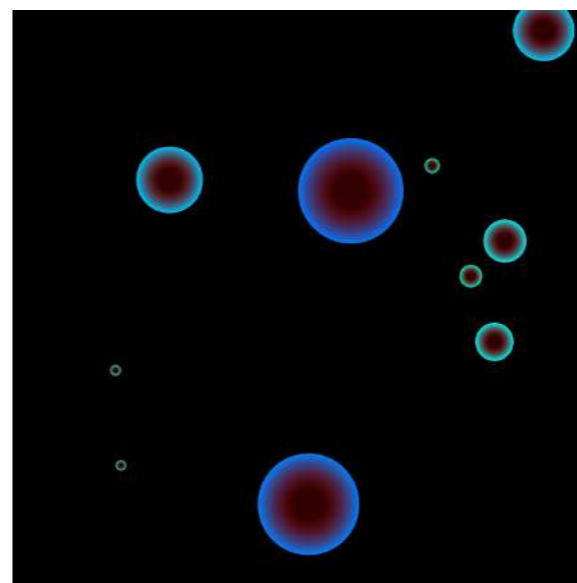
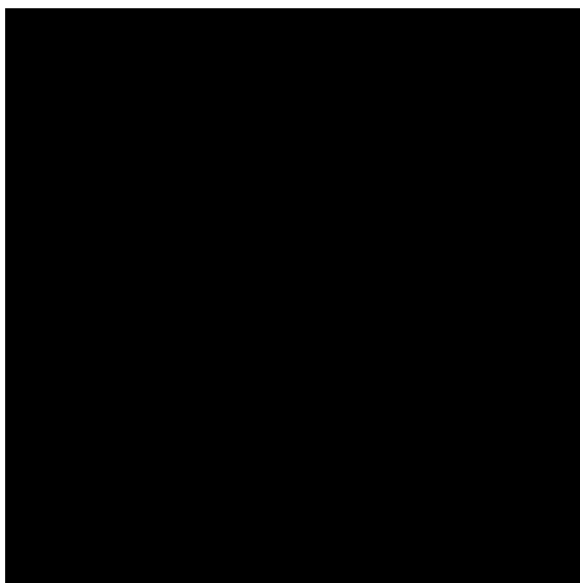
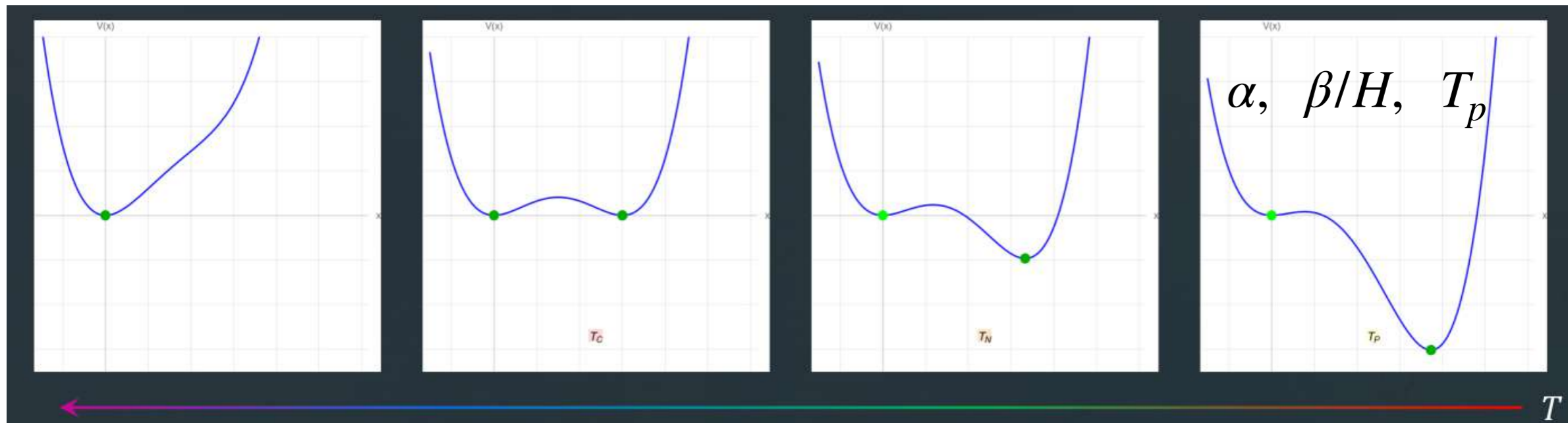
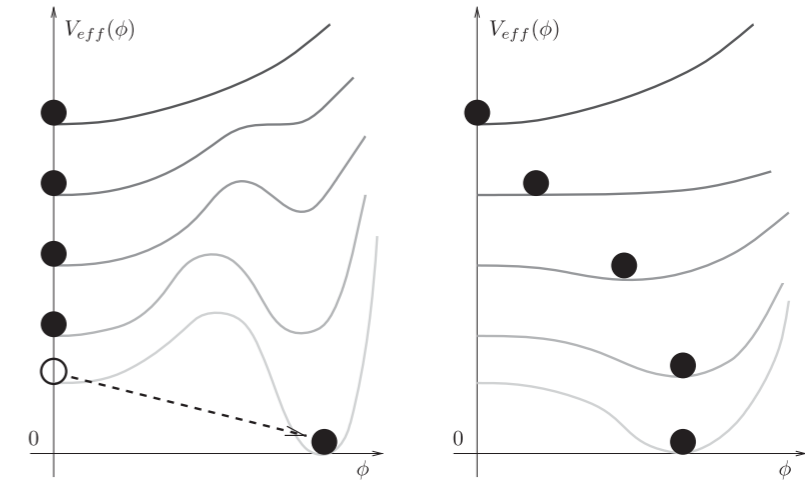
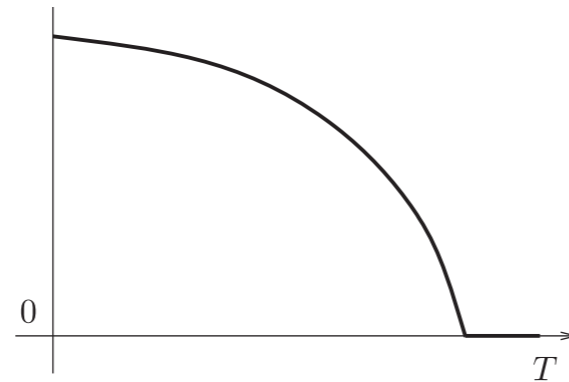
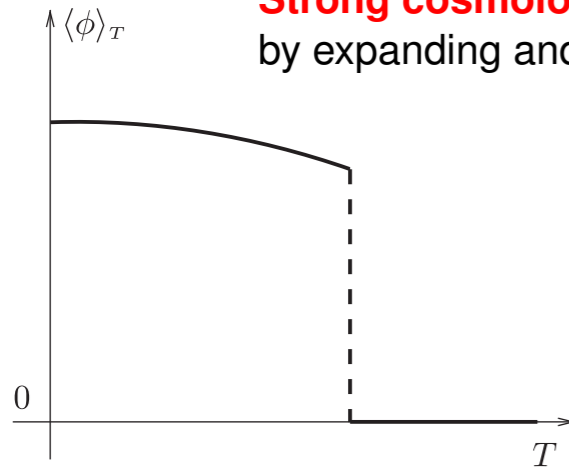
Gravitational wave detectors and sources



Wide coverage of amplitudes/frequencies in a number of GW facilities!

Phase transitions in the early Universe

Strong cosmological phase transitions (PTs) →
by expanding and colliding vacuum bubbles of new phase



Gravitational-wave power spectrum

- GW energy density per logarithmic frequency

$$h^2 \Omega_{\text{GW}} \equiv \frac{h^2}{\rho_c} \frac{d\rho_{\text{GW}}}{d \log f} \simeq h^2 \Omega_{\text{col}} + h^2 \Omega_{\text{sw}} + h^2 \Omega_{\text{MHD}}$$

Strength
 α

Inverse
duration
 β/H

Latest SGWB templates can be found in:

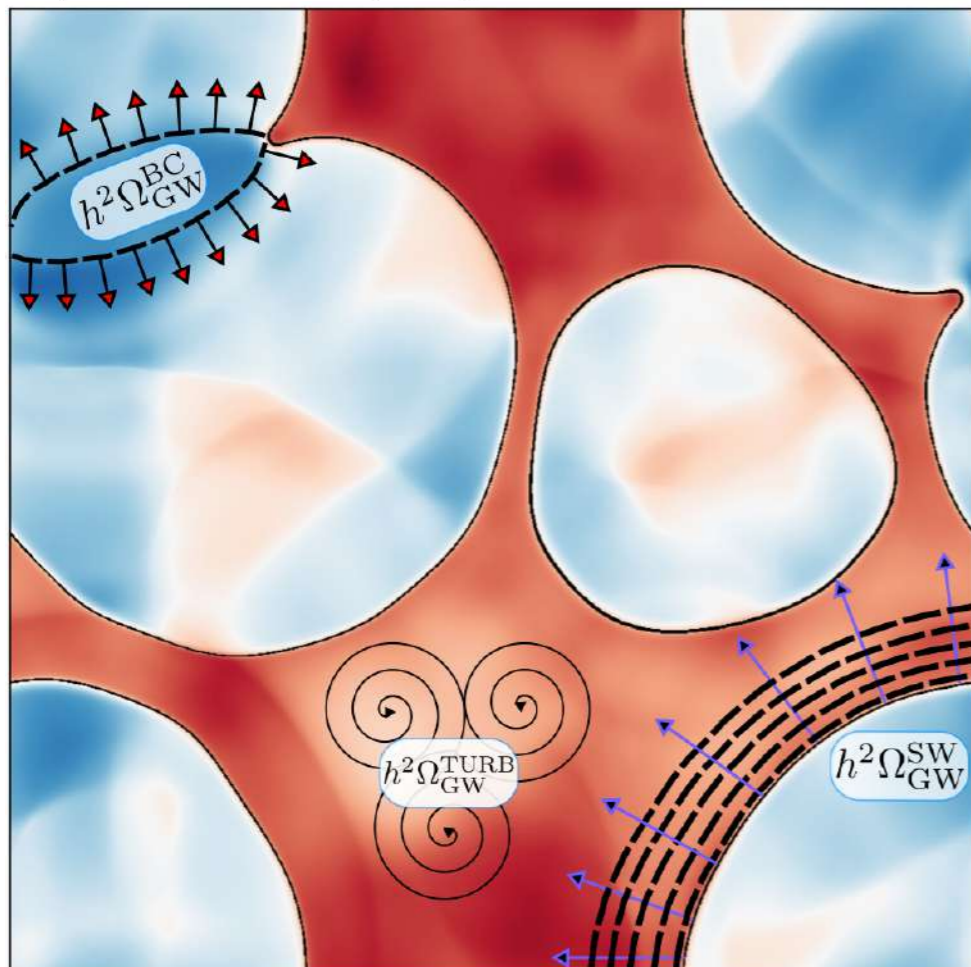
LISA Cosmology Working Group, C. Caprini et al, 2403.03723

Percolation
temperature

signal \sim amplitude \times spectral shape (f/f_{peak})

T_p

Phys.Rev.Lett. 125 (2020) 2, 021302



Primordial GWs power spectrum:

$$h^2 \Omega_{\text{GW}} = h^2 \Omega_{\text{GW}}^{\text{peak}} \left(\frac{4}{7}\right)^{-\frac{7}{2}} \left(\frac{f}{f_{\text{peak}}}\right)^3 \left[1 + \frac{3}{4} \left(\frac{f}{f_{\text{peak}}}\right)\right]^{-\frac{7}{2}}$$

Peak amplitude

Spectral function

peak frequency

$$f_{\text{peak}} \propto (\beta/H) T_*$$

$$h^2 \Omega_{\text{GW}}^{\text{peak}} \propto T_*^2 K(\alpha) f_{\text{peak}}^{-2}$$

Highly sensitive to details of FOPTs
and to underlined particle physics model

Evidence of the stochastic GW background

NANOGrav Collaboration,
Astrophys. J. Lett. 951 (2023) 1, L8; 2306.16213

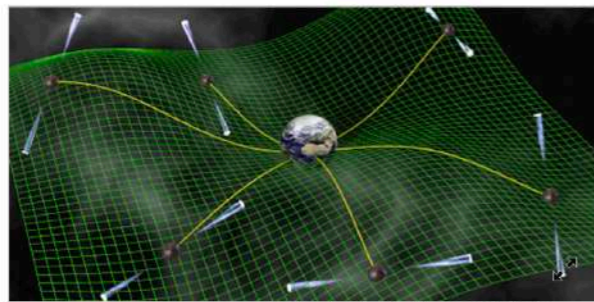


RESEARCH NEWS

Researchers Capture Gravitational-Wave Background with Pulsar “Antennae”

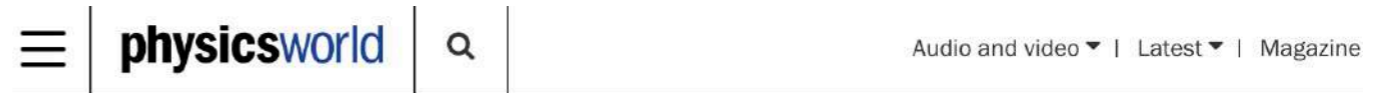
June 29, 2023 • Physics 16, 118

Four independent collaborations have spotted a background of gravitational waves that passes through our Galaxy, opening a new window on the astrophysical and cosmological processes that could produce such waves.



D. Champion/Max Planck Institute for Radio Astronomy

Pulsar timing arrays (PTAs) use a set of pulsars embedded in our Galaxy to probe the passage of gravitational waves that modulate radio signals from the pulsars. Four PTA collaborations have delivered evidence for a stochastic background of nanohertz gravitational waves.



ASTRONOMY AND SPACE | RESEARCH UPDATE

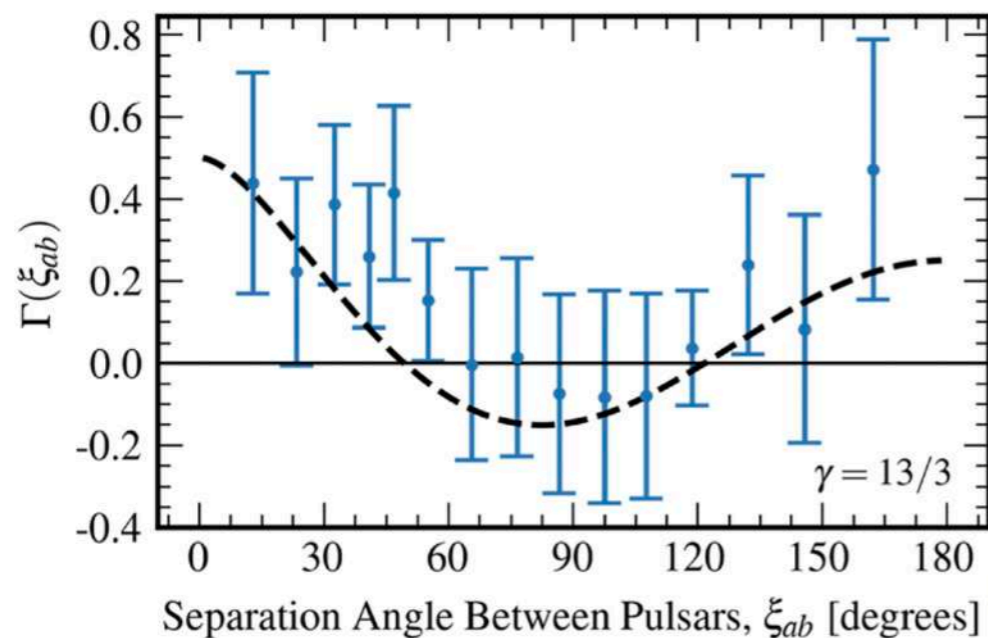
Pulsar timing irregularities reveals hidden gravitational-wave background

29 Jun 2023



Cosmic claims: researchers have used radio telescopes around the world to hunt for gravitational waves using the subtle variations in the timing of pulsars. (Courtesy: Aurore Simonnet for the NANOGrav Collaboration)

NANOGrav data vs expected correlation for SGWB



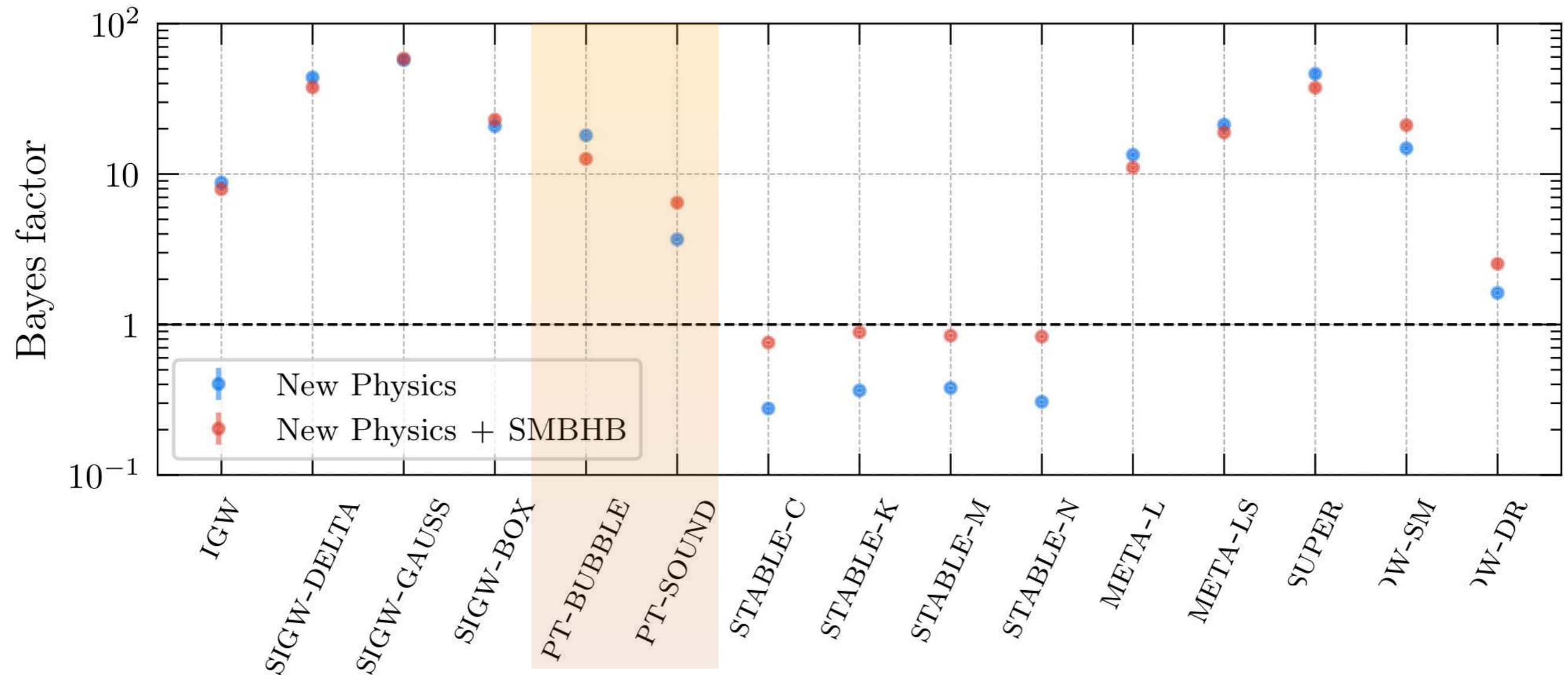
**A conservative explanation:
supermassive BH binaries**

**Can such an observation become
a window into “New Physics”?**

**The correlation between two pulsars depends on
the separation of the pulsars in the sky, measured in degrees.**

Probing New Physics at NANOGrav

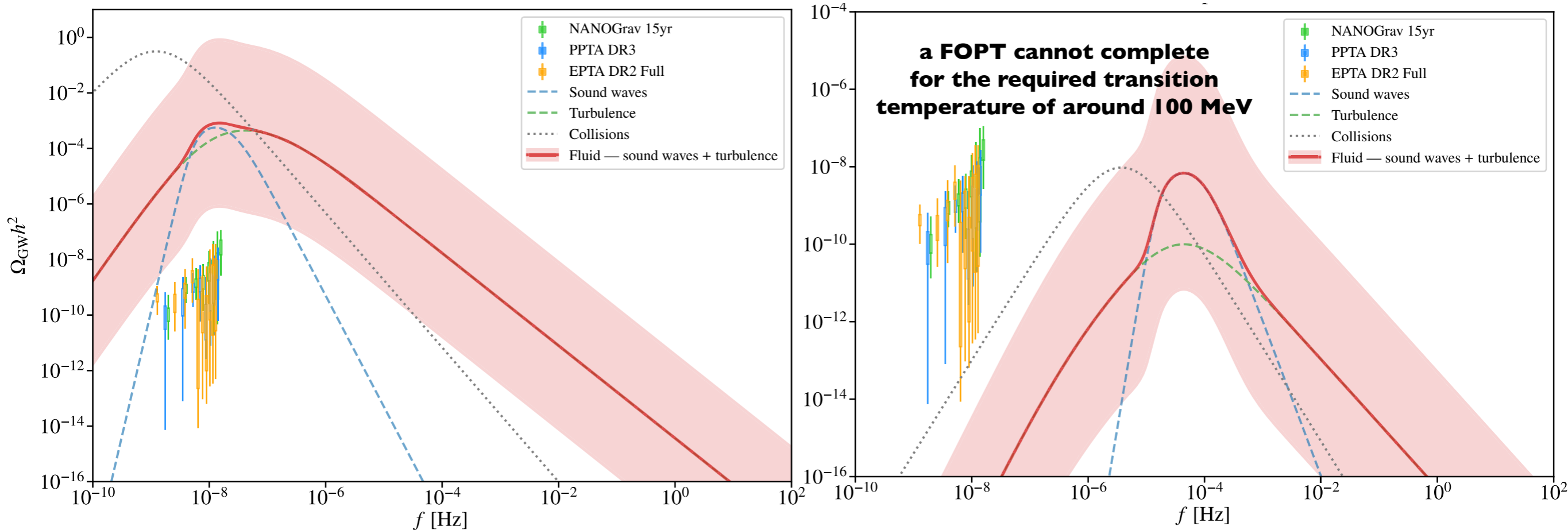
NANOGrav Collaboration,
Astrophys. J. Lett. 951 (2023) 1, L11; 2306.16219



- many models provide a better fit resulting in Bayes factors from 10 to 100
- strongly depend on modelling assumptions about the cosmic SMBHB population
- considered by many theorists as a constraint on scenarios of New Physics
- for specific models, however, making predictions requires a great care!

Call for caution in NANOGrav data interpretations

P. Athron et al, PRL 132 (2024) 22, 221001; 2306.17239



strongest supercooling for which
the FOPT is NOT completed
and percolation may not occur

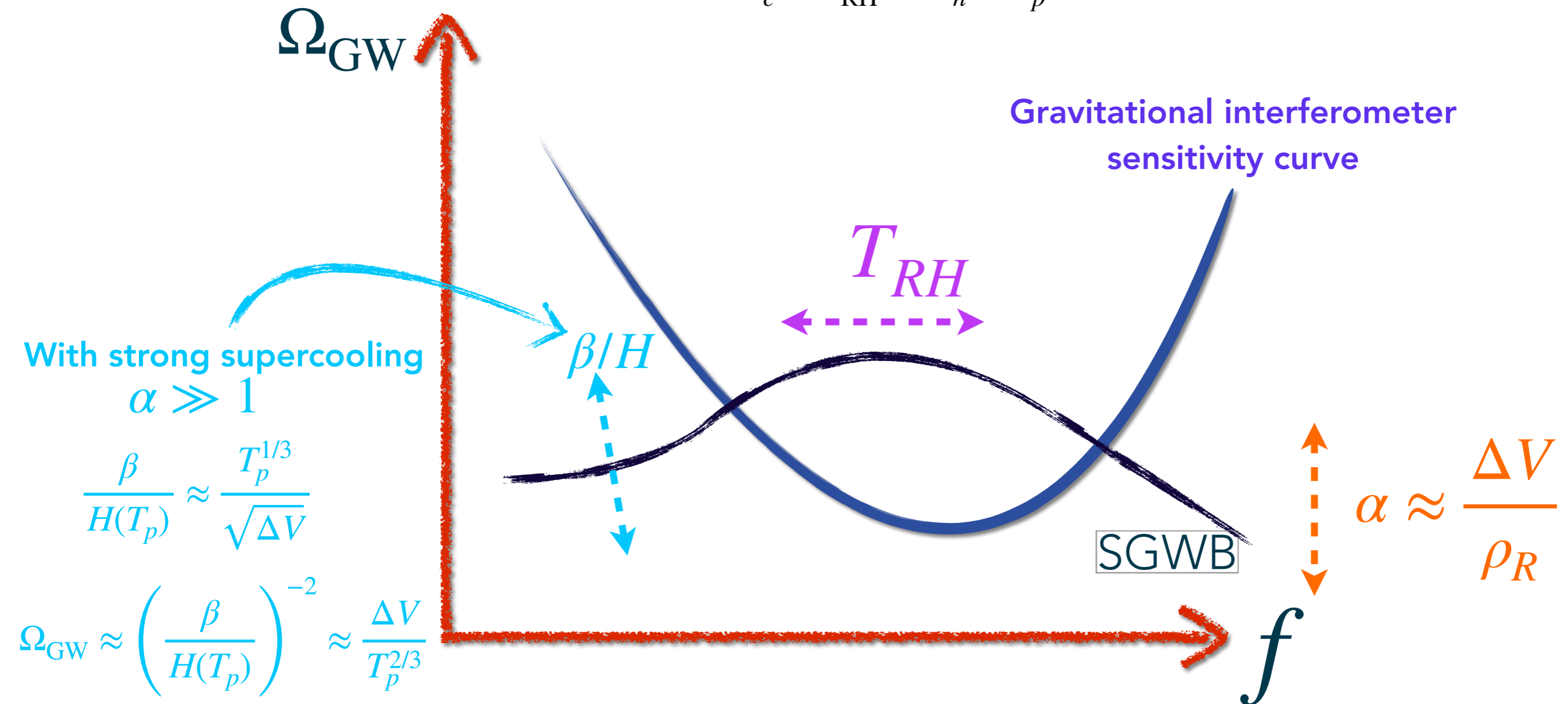
strongest supercooling for which
the FOPT is completed

period of vacuum domination hinders bubble percolation and transition completion preventing a consistent interpretation of NANOGrav signal through supercooled transitions

Stochastic GWs from supercooling

$$h^2 \Omega_{\text{GW}}^{\text{peak}} \propto \left(\frac{\alpha}{1 + \alpha} \right)^2 \left(\frac{\beta}{H(T_p)} \right)^{-2} \quad f_{\text{peak}} \propto \left(\frac{\beta}{H(T_p)} \right) \left(\frac{T_{\text{RH}}}{\text{GeV}} \right) \quad T_{\text{RH}} \propto T_p (1 + \alpha)^{1/4}$$

$$T_c > T_{\text{RH}} \gg T_n > T_p$$



Supercooling in conformal U(1)' models

generic charge assignments

| Field | U(1)' |
|---------------|---|
| Q | $\frac{1}{3}x_{\mathcal{H}} + \frac{1}{6}x_{\sigma}$ |
| u_R | $\frac{4}{3}x_{\mathcal{H}} + \frac{1}{6}x_{\sigma}$ |
| d_R | $-\frac{2}{3}x_{\mathcal{H}} + \frac{1}{6}x_{\sigma}$ |
| L | $-x_{\mathcal{H}} - \frac{1}{2}x_{\sigma}$ |
| e_R | $-2x_{\mathcal{H}} - \frac{1}{2}x_{\sigma}$ |
| \mathcal{H} | $x_{\mathcal{H}}$ |
| ν_R | $-\frac{1}{2}x_{\sigma}$ |
| σ | x_{σ} |

Tree Level CP-even scalar masses

$$M_{h_1}^{(0)} = 0 \quad M_{h_2}^{(0)} \neq 0$$

Goldstone "scalon"

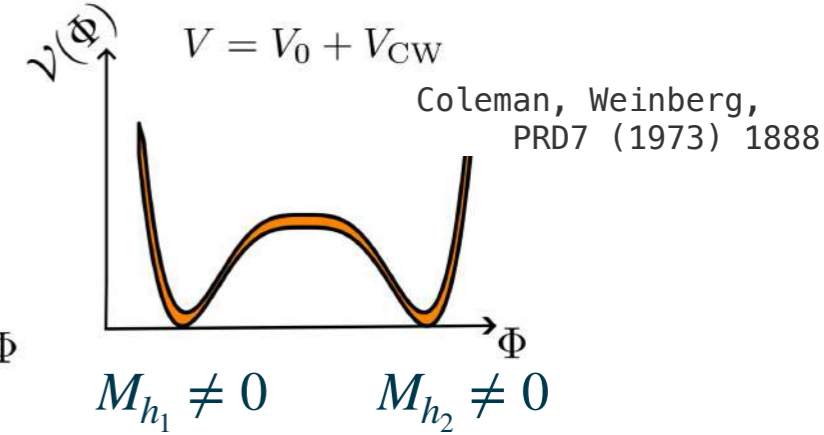
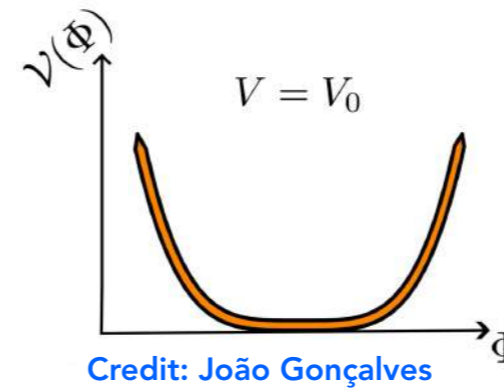
Higgs boson as pseudo-Goldstone "scalon"

Neutrino masses/mixing via Type-I seesaw mechanism

$$\mathcal{L}_{\nu} = y_{\nu}^{ij} \bar{L}_i \tilde{\mathcal{H}} \nu_{Rj} + y_{\sigma}^{ij} \bar{\nu}_{Ri}^c \nu_{Rj} \sigma + \text{h.c.}$$

scale invariant scalar potential

$$V_0(\mathcal{H}, \sigma) = \lambda_h (\mathcal{H}^\dagger \mathcal{H})^2 + \lambda_{\sigma} (\sigma^\dagger \sigma)^2 + \lambda_{\sigma h} (\mathcal{H}^\dagger \mathcal{H})(\sigma^\dagger \sigma)$$

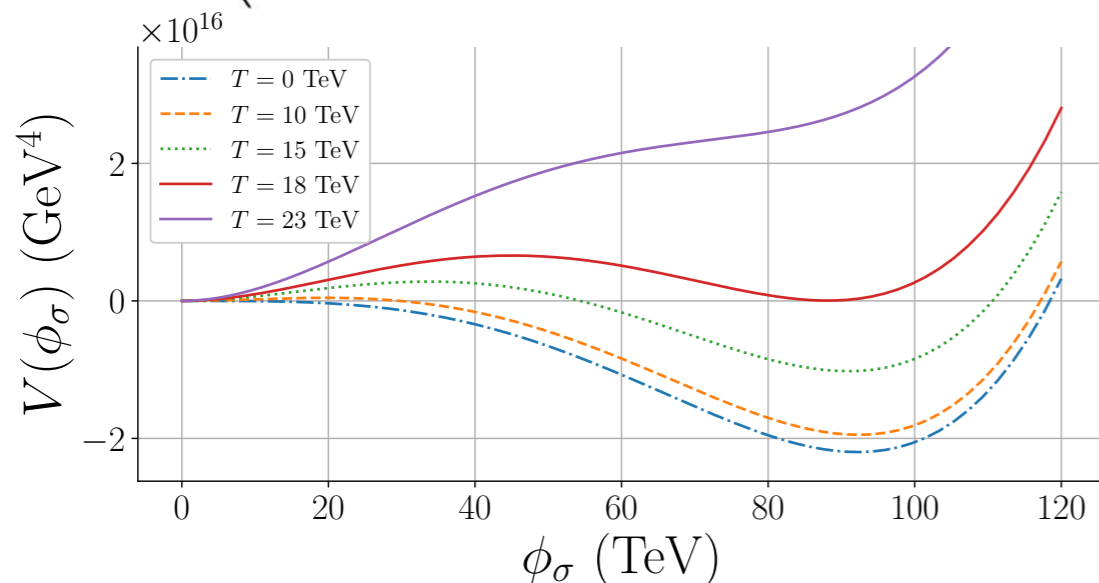


[Gildener, Weinberg, PRD13 (1976) 3333]

High-temperature EFT potential

$$V_{\text{eff}}^{\text{HT}} = \phi_{\sigma}^4 \left(-\frac{g_L^4}{2\pi^2} - \frac{g_L^3}{2\sqrt{2}\pi} + \frac{\lambda_{\sigma}}{4} + \frac{\ln 2 \left(\left[\sum_{i=1}^3 [y_{\sigma}^4]_{ii} \right] \right)}{32\pi^2} \right) - \phi_{\sigma}^3 \frac{4g_L^3 T}{3\pi} + \phi_{\sigma}^2 \left(\frac{g_L^2 T^2}{2} - \frac{g_L^3 T^2}{\sqrt{2}\pi} + \frac{T^2}{48} \sum_{i=1}^3 [y_{\sigma}^2]_{ii} \right)$$

negative cubic term is generated!



⇒ dynamical symmetry breaking

⇒ potential barrier persists till very low T

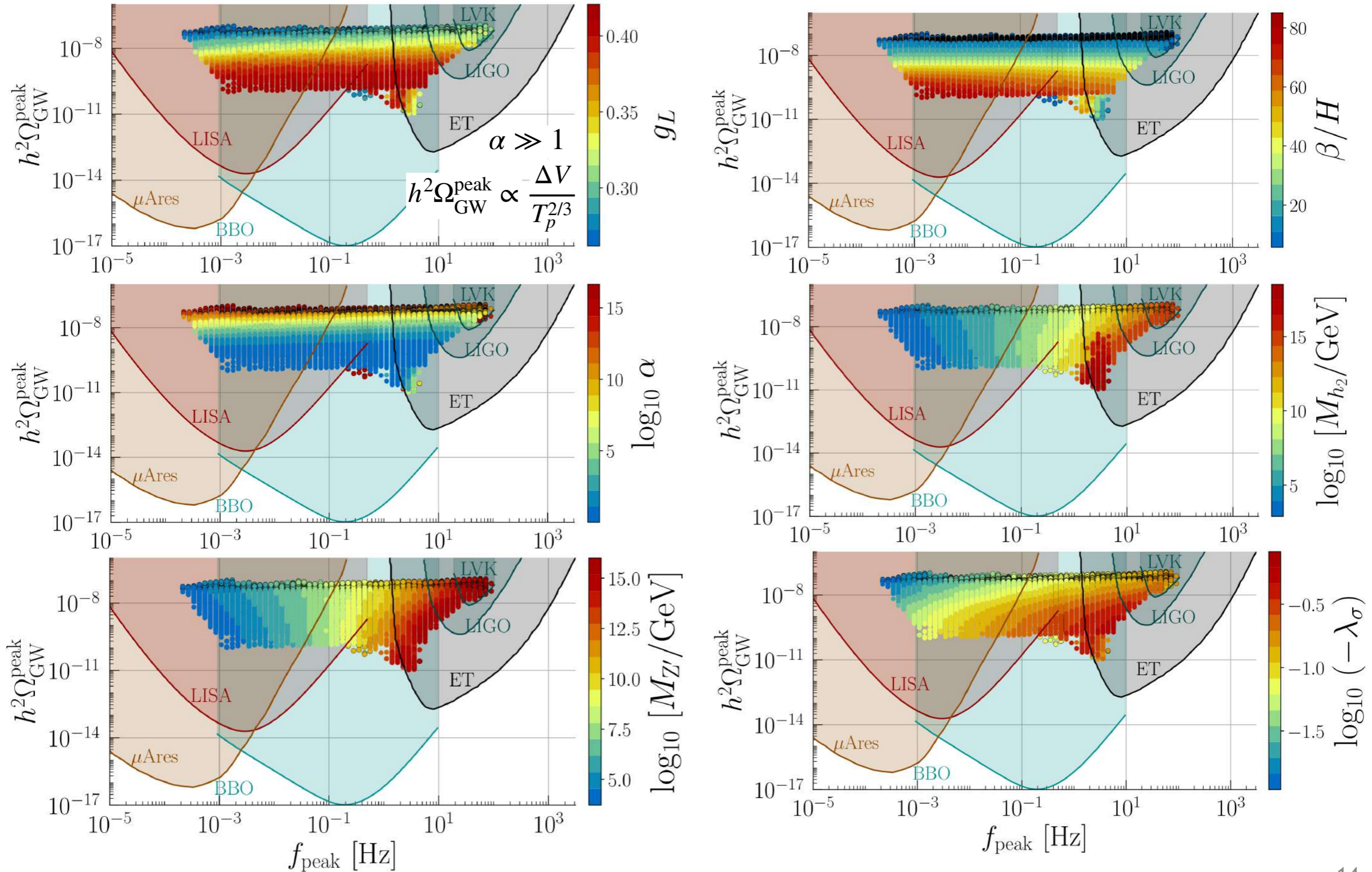
⇒ long-lasting FOPT $\beta/H \sim \mathcal{O}(10 - 100)$

⇒ ΔV is maximised leading to $\alpha \approx \frac{\Delta V}{\rho_R} \gg 1$

SGWB predictions in conformal B-L theory

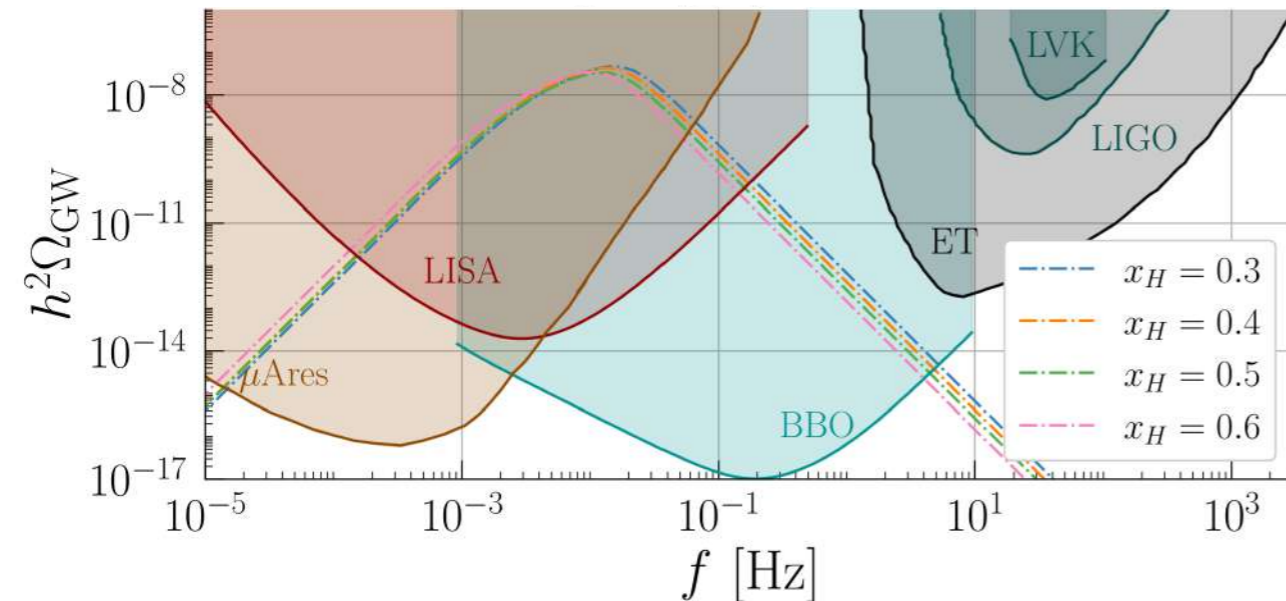
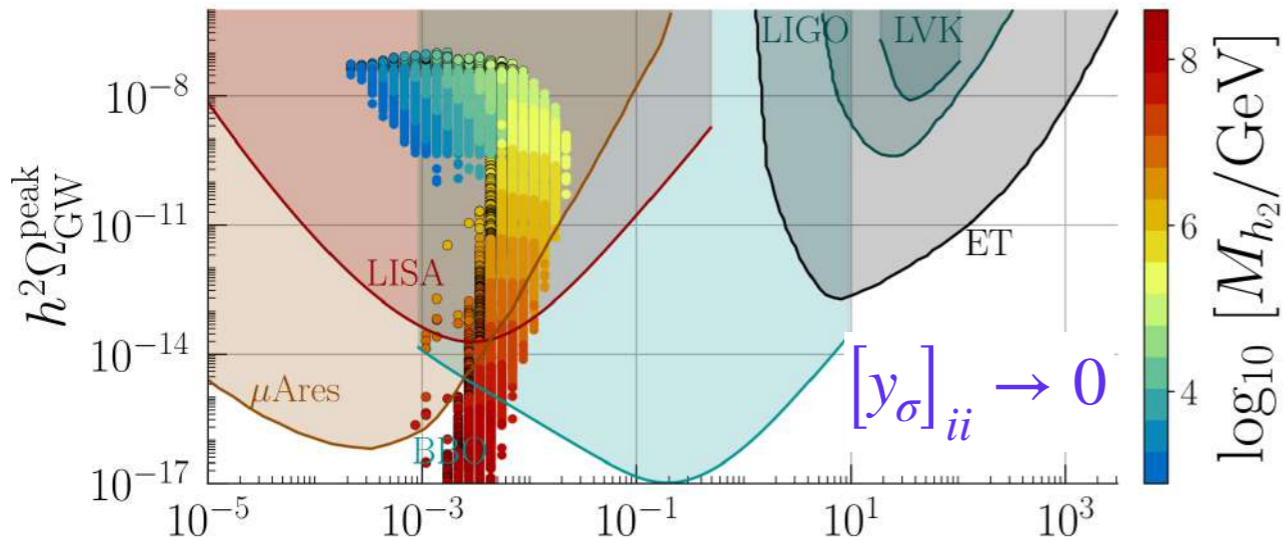
$x_\sigma = 2$ and $x_H = 0$

$\beta/H \gtrsim 8$ from PBH constraints [Y. Gouttenoire, T. Volanski, 2305.04942]

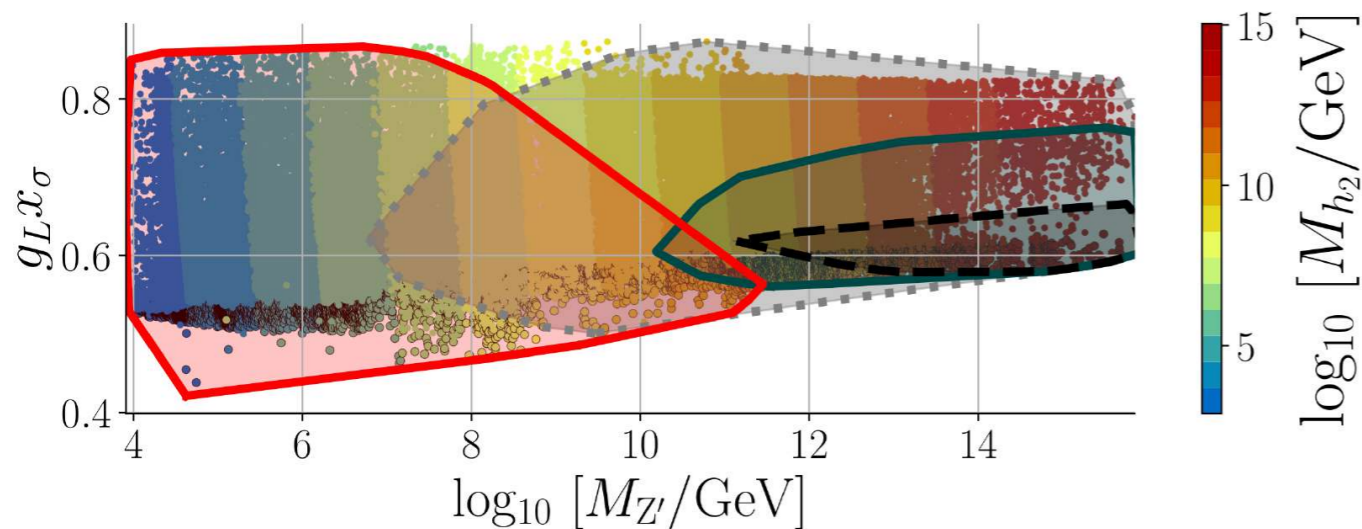


SGWB predictions: role of neutrinos and charges

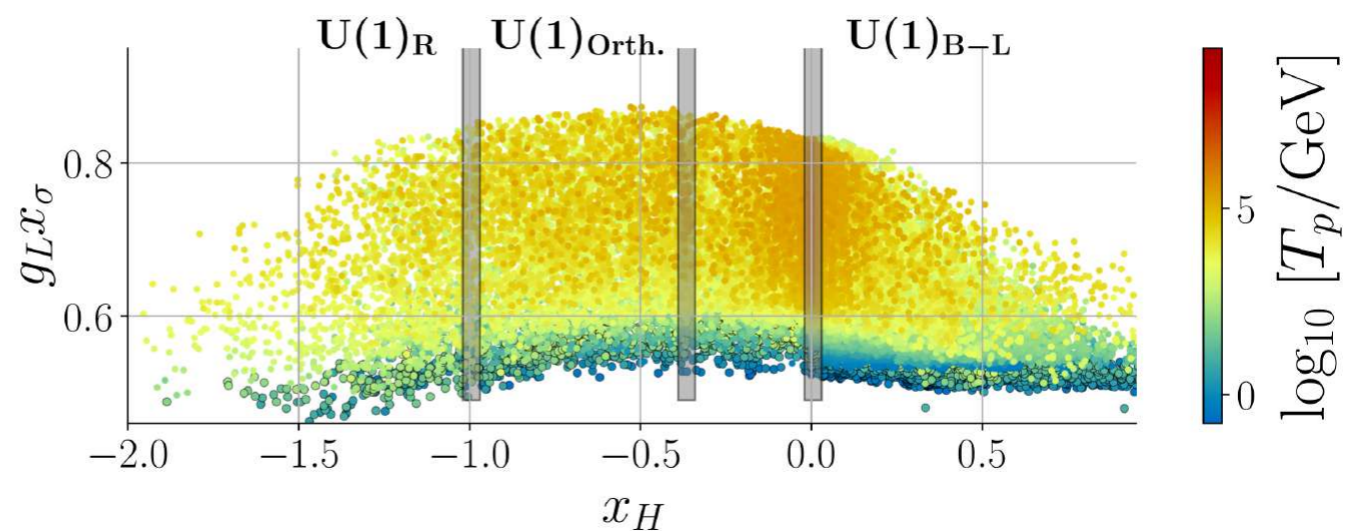
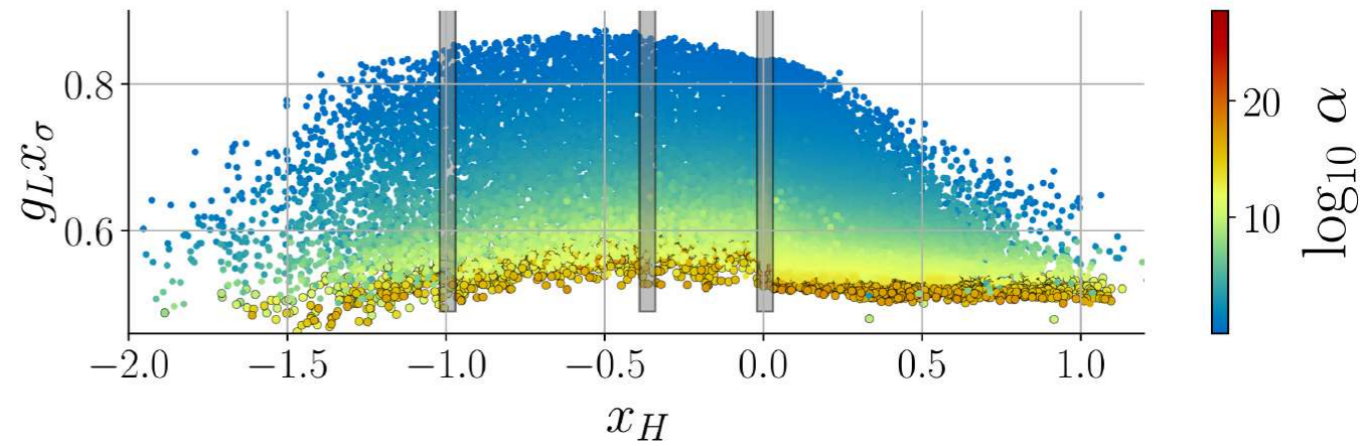
LIGO/ET signal \Leftrightarrow neutrino mass generation



SNR > 10 for observable SGWB



$(-1, 2)$ $(-\frac{16}{41}, 2)$ $(0, 2)$ (x_H, x_σ)
 $U(1)_R$ $U(1)_{Orth.}$ $U(1)_{B-L}$



LIGO+Virgo+KAGRA (LVK) excluded:

$$10^{12} \text{ GeV} < M_{h_2} \sim M_{Z'} < 10^{16} \text{ GeV}$$

$$g_L x_\sigma \sim 0.6$$

- LISA ———
- ET (dotted)
- LIGO-O5 ——— (green)
- LVK - - - - - (dashed)

LISA+ET+LIGO will probe:

$$M_{h_2} > 1 \text{ TeV}, M_{Z'} > 10 \text{ TeV}$$

$$0.55 \lesssim g_L x_\sigma \lesssim 0.8$$

Conclusions

- ◆ Vacuum domination hinders bubble percolation in the frequency domain of PTAs making it difficult to interpret NANOGrav data due to supercooled phase transitions
- ◆ SGWB detection at LISA+ET+LIGO probes supercooling at higher temperatures (hence, larger frequencies) setting bounds on parameters of conformal models with new gauge symmetries
- ◆ Simultaneous explanation of neutrino masses in the $U(1)'$ conformal model is crucial for high-frequency SGWB searches
- ◆ LISA+ET+LIGO can either rule out most of the parameter space challenging the hypothesis of supercooling and scale invariance, or lead to a ground-breaking discovery